



## TiO<sub>2</sub> films with rich bulk oxygen vacancies prepared by electrospinning for dye-sensitized solar cells

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### HIGHLIGHTS

- ▶ Transparent TiO<sub>2</sub> films were prepared by electrospinning.
- ▶ TiO<sub>2</sub> films with rich bulk oxygen vacancies induce large  $V_{oc}$  and FF in DSCs.
- ▶ We first used an unsensitized cell to investigate the trap-to-trap electron transport.
- ▶ The results from PL measurement show useful information.

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### ABSTRACT

Highly transparent nanocrystalline TiO<sub>2</sub> films have been fabricated by electrospinning (ES) technique based on a transmutation process from as-spun nanofibers with an appropriate amount of triethanolamine (TEOA) added to the precursor. A possible evolution mechanism of the transparent nanocrystalline TiO<sub>2</sub> films is proposed. It is found that the films prepared via transmutation from electrospun nanofibers possess rich bulk oxygen vacancies (BOVs, PL band at 621–640 nm) by using photoluminescence (PL) spectroscopy. Contrastively, the dominant peak in PL spectrum of the spin-coated film is the emission from surface oxygen vacancies (SOVs, PL band at 537–555 nm). The electrospun TiO<sub>2</sub> films with rich BOVs induce large open-circuit voltage ( $V_{oc}$ ) and fill factor (FF) improvements in dye-sensitized solar cells (DSCs), and thus a large improvement of energy conversion efficiency ( $\eta$ ). In addition, these performance advantages are maintained for a double-layer cell with a doctor-bladed  $\sim 7 \mu\text{m}$  top layer (P25 nanometer TiO<sub>2</sub>, Degussa) and an electrospun  $\sim 3 \mu\text{m}$  bottom layer. The double-layer cell yields a high  $\eta$  of 6.01%, which has increased by 14% as compared with that obtained from a 10  $\mu\text{m}$  thick P25 film.

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### 1. Introduction

Dye-sensitized solar cells (DSCs) have attracted extensive interest in the past decade due to the prospects of low-cost, environment-friendly, high-performance and of new applications such as transparent solar cells and flexible solar cells [1–4]. Nanostructured metal oxides are one of the key factors in determining the photovoltaic-characteristics of DSCs, because the nanostructured networks will act as a scaffold for the dyes and a transfer media for photogenerated electrons to the transparent conductive oxide layer. Therefore, the properties of the metal oxides, such as

the defects in the materials, must greatly influence the photovoltaic-characteristics of DSCs.

In DSCs, the photogenerated electrons would transfer directly from the conduction band or the surface traps of TiO<sub>2</sub> nanoparticles [5–7] to acceptors in electrolyte. Thus the traps located at the TiO<sub>2</sub> surface would lower the electron lifetime in DSCs and resulting in poor cell performance [8]. For this reason, various methods have been developed to improve the cell performance by minimizing the possible recombination pathways occurring at the TiO<sub>2</sub>/dye/electrolyte interface [9–11]. However, the electron traps in the bulk may behave differently. As the trap sites are filled, the potential difference between the highest filled traps and the conduction band decreases. This may increase excitation possibility of the trapped electrons from the traps to the conduction band. Furthermore, the bulk traps trap and release electrons only

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with the conduction band [6], namely the bulk traps are unable to transfer electrons to the acceptor species in electrolyte solution. Nelson et al. [7] demonstrated that multiple trapping with a broad energetic distribution of electron traps is responsible for the slow recombination kinetics in DSCs. Thus, using TiO<sub>2</sub> films with rich bulk electron traps may be a practical strategy to achieve the high photovoltaic performance, which is of great importance but has not been done.

As a simple and cost-effective way to prepare semiconductor oxides with 1D nanostructure, electrospinning technique has attracted extensive interest in various areas, including dye-sensitized solar cells [12], gas sensor [13], photo-catalysis [14], lithium-ion batteries [15], and transparent conductive films [16]. Compared with other methods, electrospinning possesses many advantages, such as large-area preparation, on both conductive and insulating substrates, and cost-effective. It is reported that different nanostructures, for example, nanofibers [17], nanotubes [18] and nanobelts [19] have been fabricated by electrospinning technique. Here, we report the preparation of highly transparent nanocrystalline TiO<sub>2</sub> films by electrospinning. The resultant films possess rich bulk oxygen vacancies and can induce large open-circuit voltage ( $V_{oc}$ ) and fill factor (FF) improvements in DSCs.

## 2. Experimental details

### 2.1. Preparation of the transparent nanocrystalline TiO<sub>2</sub> films

To prepare the precursor solutions, 0.5 g tetrabutyl titanate ( $\geq 98\%$ ) was dissolved in a mixture of 1 ml acetic acid and 1 ml ethanol. The obtained solution was stirred for 10 min to form a sol. Then, a solution containing 0.19 g poly (vinyl pyrrolidone) (PVP, Sigma Aldrich, Mw  $\approx 1,300,000$ ), 1.5 ml ethanol and 0.5 ml N,N-Dimethylformamide (DMF) were mixed with the above sol. After stirring the obtained solution for 20 min, tri-ethanolamine (TEOA) was added. Finally, the precursor was obtained after stirring the solution for 30 min. The applied voltage and the distance between the nozzle and the collector were fixed at 6.5 kV and 10 cm, respectively. To achieve the transmutation of the electrospun nanofibers, the humidity level was controlled at 50–55%.

### 2.2. Assembling of DSCs

All TiO<sub>2</sub> electrodes were sintered at 500 °C for 2 h. For the post-treatment, the TiO<sub>2</sub> films were dipped into a 100 mM TiCl<sub>4</sub> aqueous solution for 20 min at 70 °C, then sintered again at 500 °C for 30 min. At 80 °C in the cooling, the sintered TiO<sub>2</sub> electrodes were immersed into a 0.3 mM solution of N-719 dye in a mixture of tert-butyl alcohol and acetonitrile (volume ratio of 1:1) and kept in the dark at room temperature for 12 h to complete the adsorption of the dye. The platinized counter electrode was fabricated by spin-coating a 6.0 mM isopropanol solution of H<sub>2</sub>PtCl<sub>6</sub>•6H<sub>2</sub>O, then annealing at 400 °C for 20 min. The electrolyte was a solution of 0.1 M LiI, 0.6 M 1,2-dimethyl-3-propylimidazorium iodide, 0.06 M I<sub>2</sub> and 0.5 M 4-tert-butylpyridine in acetonitrile. The preparation of a P25 based paste for reference is described in our previous work [20]. The cell area was fixed at 0.45 cm<sup>2</sup>.

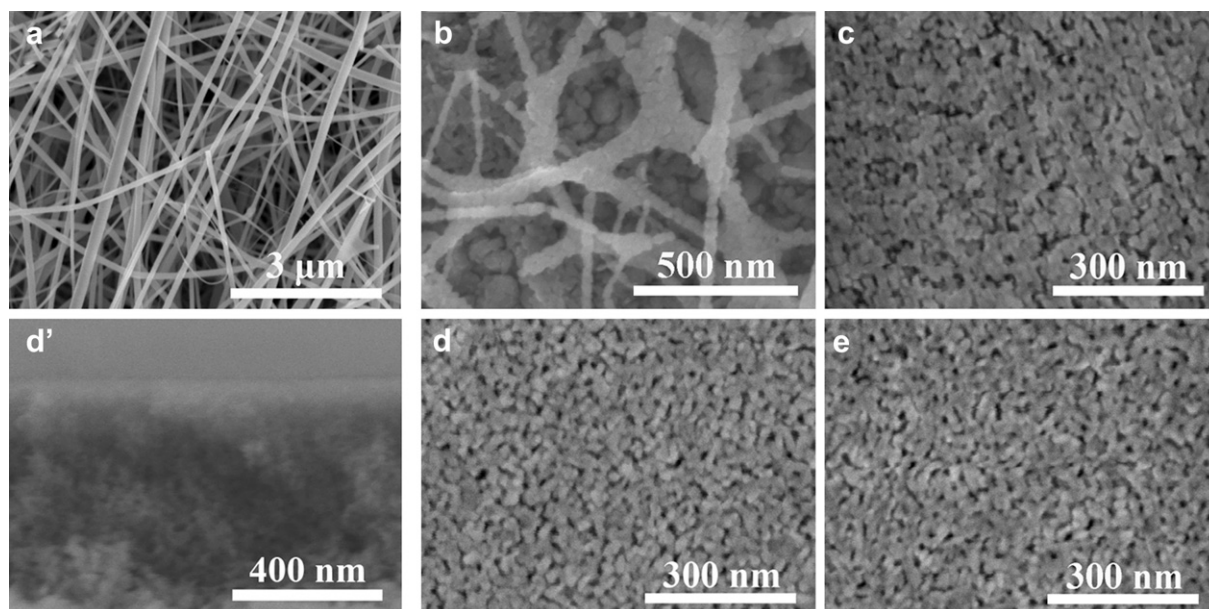
### 2.3. Characterization

Morphologies and optical properties were studied by field emission scanning electron microscopy (FESEM, Hitachi, S-4800) and UV–VIS spectrophotometer (TU-1901), respectively. X-ray diffraction (XRD) (Philips X'pert Pro MPD (Cu K $\alpha$  ray)) was used to characterize the crystallographic property. PL measurements were carried out on a Raman spectroscopy (JY-HR800) with a 325 nm line of 50 mW diode-pumped solid state laser at room temperature. AC impedance spectroscopy was performed on an electrochemical system (RST-5200, Zhengzhou Shiruisi Instrument Technology Co., Ltd, China) with frequency range from 0.1 to 10<sup>5</sup> Hz and the oscillation amplitude of 10 mV. Thicknesses of the TiO<sub>2</sub> films were obtained by using a surface profile measurement system (Veeco, Dektak8).

## 3. Results and discussion

### 3.1. The effect of TEOA on TiO<sub>2</sub> films

Without TEOA in the precursor, a nanofiber mat (Fig. 1(a)) was formed, which is a common result achieved by many other groups



**Fig. 1.** FESEM images of the electrospun samples with precursor solutions consisting of 0 (a), 3.8 (b), 5.0 (c), 6.3 (d) and 8.5 wt.% (e) TEOA. Cross-sectional SEM image of the film with 6.3 wt.% TEOA (d'). All the samples were annealed at 500 °C for 2 h.

[21]. With a small amount of TEOA added, such as 3.8 wt.%, a nanofibers-covered nanocrystalline film is obtained (Fig. 1(b)). This film is consisted of hierarchical aggregates with average diameter ranging from 65.3 nm to 135.5 nm, and the size of the primary nanoparticles in the aggregates is between 14.9 nm and 40.1 nm. As the proportion of TEOA increased, the nanofibers disappeared gradually and films were formed with an increasing uniformity of the particle sizes (Fig. 1(c)–(e)). Moreover, when the proportion of TEOA was increased from 3.8 wt.% to 8.5 wt.%, the average particle sizes of the films decrease from 32.0 nm to 18.5 nm (Fig. 2). Cross-sectional SEM image of the films with 6.3 wt.% TEOA shown in Fig. 1(d') reveals that the film consists of uniform nanoparticles with an average size of  $\sim 23$  nm, which is consistent with the surface morphology shown in Fig. 1(d). In addition, the films exhibit high transparency and large area uniformity as shown in Fig. S1.

### 3.2. Formation mechanism of the transparent nanocrystalline TiO<sub>2</sub> films

A possible evolution mechanism of the transparent nanocrystalline TiO<sub>2</sub> films is shown in Fig. 3. With TEOA in the precursor, the nanofibers cannot keep their shape and may transmute from nanofiber mat to a continuous gel layer due to the intrinsic hygroscopic and viscous properties of TEOA. Because it is viscous, the TEOA is hardly to evaporate and make the nanofibers deform easily. Because it is hygroscopic, TEOA can absorb the moisture in the air which will redound to the nanofibers deformation process. Here, a high humidity of 50–55% is required. It was reported that TEOA acts as an inhibitor of the coagulation of the growing hydroxide particles [22]. With a small amount of TEOA was added, such as 3.8 wt.%, the transmutation process was not complete and a loose gel layer may form. During the process, coagulation of the growing hydroxide particles happened due to the high humidity and low content of TEOA, resulting in TiO<sub>2</sub> particles with larger size. While with a large amount of TEOA was added, such as 8.5 wt.%, the transmutation process was complete, resulting in a compact gel layer. Because of the high concentration of TEOA, coagulation of the growing hydroxide particles is hindered, leading to TiO<sub>2</sub> particles with smaller particle sizes and better uniformity.

To assess the validity of the above mechanism, samples were fabricated by spin-coating by using the same precursors as that in

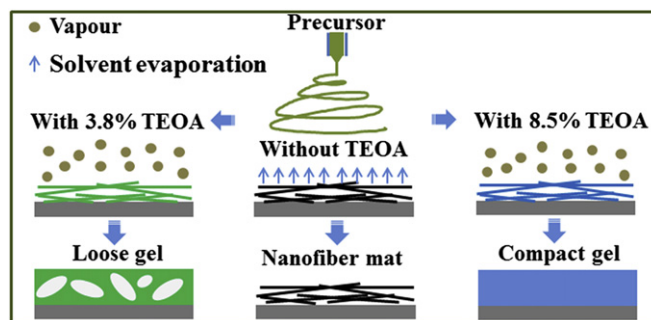


Fig. 3. Schematic diagram showing the effect of TEOA on the formation of the transparent nanocrystalline TiO<sub>2</sub> films.

electrospinning. As shown in Fig. 2, the particle sizes of the spin-coated samples stay the same which was in sharp contrast to that of electrospinning samples. This is because the gel layer formed on the substrate by spin-coating method is compact and is irrelevant to the proportion of TEOA. While for electrospinning, the transmutation happens on the substrate under the combined action of TEOA and humidity, which is dependent on the proportion of TEOA.

### 3.3. Structural characterization

Fig. 4 shows the XRD spectra of TiO<sub>2</sub> nanocrystalline films prepared by electrospinning and spin-coating using the precursor with 8.5 wt.% TEOA. The films were annealed in air at 500 °C for 2 h. Both samples show a typical anatase phase. PL measurement is a powerful technique to investigate the electronic structure and energy levels in the band-gap of semiconductor materials, which could influence the charge transport and recombination property in DSCs. Herein, we investigated the room-temperature PL spectra of the samples fabricated with various TEOA concentrations by electrospinning (Fig. 5(a)–(c)). For comparison, PL spectrum of the sample prepared by spin-coating with 8.5 wt.% TEOA (SC-8.5) is also shown in Fig. 5(d). It is necessary to indicate that the particle size of the film prepared by spin-coating was controlled to be  $\sim 17.8$  nm by changing spin-coating parameters (Fig. S2), which is comparable to the particle size of the film by electrospinning with 8.5 wt.% TEOA (ES-8.5, Fig. 2). As shown in Fig. 5(a)–(d), all the PL bands can be

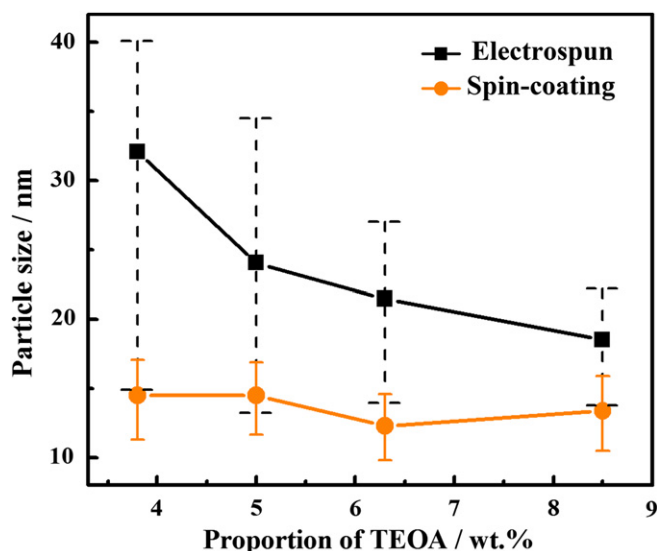


Fig. 2. Particle sizes as function of the proportion of TEOA.

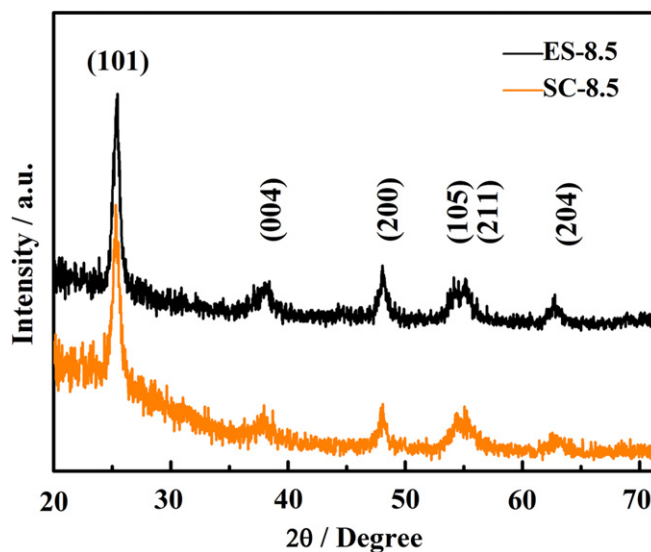
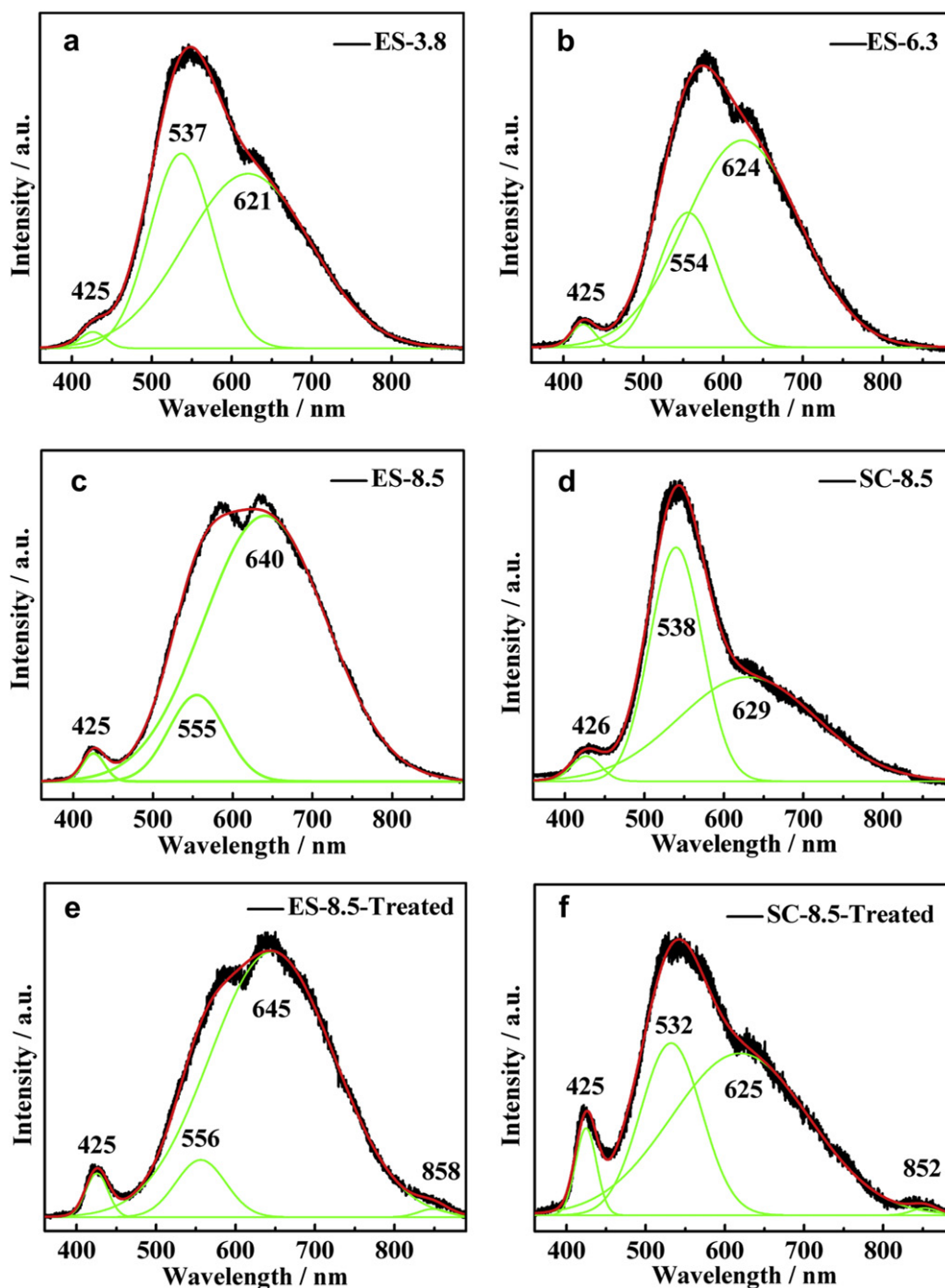


Fig. 4. XRD spectra of TiO<sub>2</sub> nanocrystalline films prepared by electrospinning and spin-coating using the precursors with 8.5 wt.% TEOA.



**Fig. 5.** PL spectra of nanocrystalline TiO<sub>2</sub> films prepared by electrospinning with different proportions of TEOA: 3.8 (a), 6.3 (b) and 8.5 wt.% (c), and by spin-coating with 8.5 wt.% TEOA (d). PL spectra of the films prepared by electrospinning (e) and spin-coating (f) with 8.5 wt.% TEOA after treatment by TiCl<sub>4</sub> aqueous solution.

fitted into three Gaussian peaks with wavelengths of 425–426 nm, 537–555 nm and 621–640 nm. As the proportion of TEOA increases, the intensity of the peaks at 621–640 nm increases and becomes dominant when the proportion of TEOA reaches 8.5 wt.%, which is in sharp contrast to that of the samples prepared by spin-coating method (Fig. 5(d)). To achieve successful transmutation of the electrospun nanofibers hygroscopic TEOA was added to the precursor and a high humidity level of 50–55% was kept. Thus, the significant difference in PL measurement between the two samples may result from the quick hydrolyze of tetrabutyl titanate during electrospinning.

Generally, PL spectra of anatase TiO<sub>2</sub> materials are interpreted as the emission from three physical origins: self-trapped excitons

[23,24], oxygen vacancies (OVs) [25,26] and surface states [27]. In anatase TiO<sub>2</sub> materials, lattice relaxation energy ( $E_{lr}$ ) is larger than the energy reduction ( $B = \nu|y|$ ) in the delocalization process by transferring from one site to a neighbor site [23]. Under this condition, the self-trapped excitons on the TiO<sub>6</sub> octahedra recombine with impurity-trapped excitons radiatively. As shown in Fig. 5(a)–(d), the 425 nm bands hold constant regardless of the proportion of TEOA and the preparation method. Moreover, the peak position of the 425 nm band is very close to that in Ref. [24] (425 nm) and Ref. [26] (412 nm). So the 425 nm band should be accordingly assigned to self-trapped excitons (STE) localized on TiO<sub>6</sub> octahedra. PL bands at the long wavelength side of anatase TiO<sub>2</sub> nanoparticles have been attributed to the OVs [24–26].

Shallow traps identified with OVVs were established at 0.8 eV below the conduction band [26]. Forss and Schubnell [27] have attributed the band corresponding to shallow traps to surface states. In our case, the 537–555 nm (0.89–0.96 eV) bands in Fig. 5(a)–(d) are very close to that in Ref. [26] (0.80 eV) and Ref. [24] (0.82 eV). Therefore, we assign the PL bands at 537–555 nm to OV located at the surface of the TiO<sub>2</sub> nanoparticles (SOVs), which can be confirmed later. There are another energy levels resulting from OV in TiO<sub>2</sub> at about 1.18 eV (614 nm) [28] below conduction band. Accordingly, we assign the PL bands at 621–640 nm to emission from oxygen vacancies. Furthermore, these oxygen vacancies are in the bulk of TiO<sub>2</sub> nanoparticles.

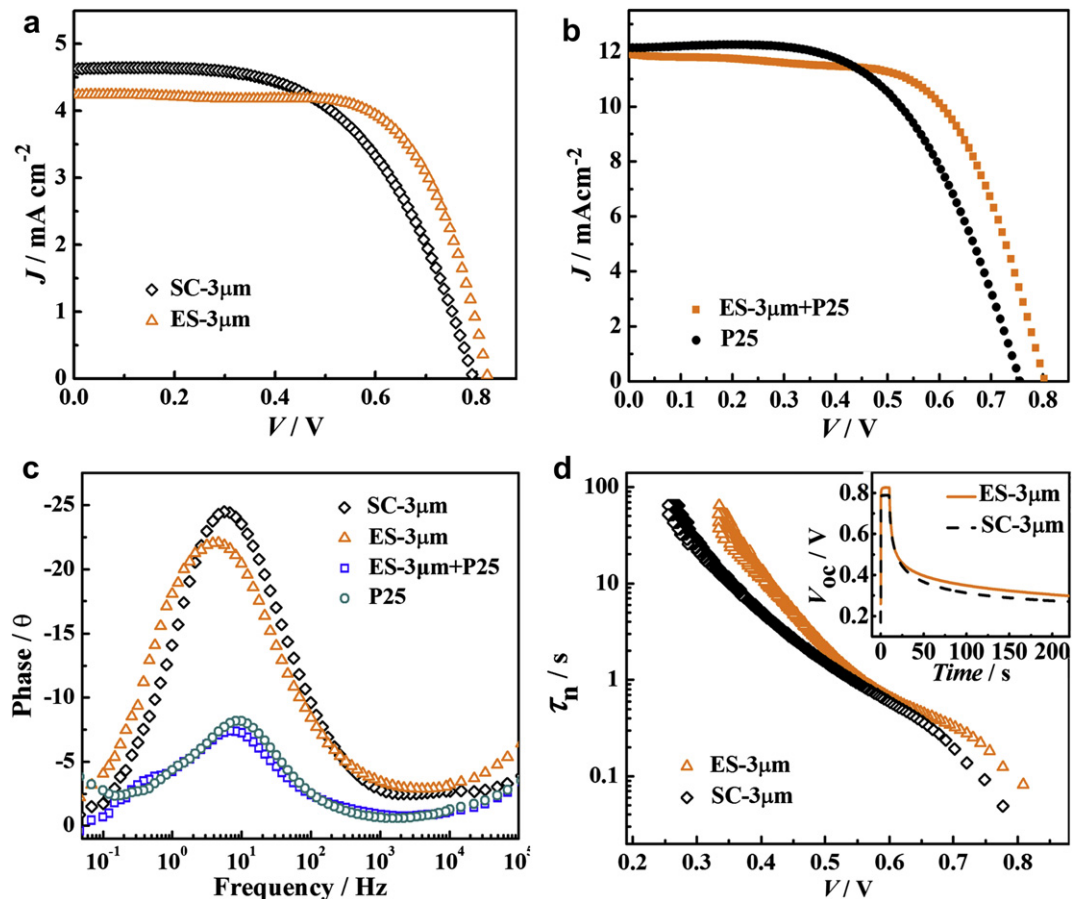
To demonstrate the origin of the PL bands, a rutile TiO<sub>2</sub> passivation layer was formed on the surface of the nanoparticles by a 100 mM TiCl<sub>4</sub> aqueous solution treatment at 70 °C for 20 min. In rutile TiO<sub>2</sub> materials,  $B > E_{\text{tr}}$ , the excitons are stabilized in the Free State. In this case the PL spectrum of rutile TiO<sub>2</sub> materials is either resulted from free exciton emission at low-temperature, or infrared emission from its intrinsic properties [23,29,30]. As shown in Fig. 5(e) and (f), the PL band at 852–858 nm proved the existence of rutile TiO<sub>2</sub> coating. No peak from free exciton emission of rutile TiO<sub>2</sub> is observed. It is found that the intensity of 537–555 nm bands decreases sharply after surface passivation, while the intensity of 425 nm and 621–640 nm bands are almost unchanged after treatment. Therefore, the 537–555 nm and 621–640 nm bands originate from the defects located at surface and in bulk of nanoparticles, respectively. In addition, the two bands are dominant peaks of the corresponding samples. These provide convenience to

investigate the effect of different defects on the photovoltaic properties of DSCs.

#### 3.4. Influence on photovoltaic-characteristics of DSCs

$J$ – $V$  curves recorded from the cells with 3  $\mu\text{m}$  thick ES-8.5 (ES-3  $\mu\text{m}$ ) and SC-8.5 (SC-3  $\mu\text{m}$ ) TiO<sub>2</sub> films as photoanodes are shown in Fig. 6(a). ES-3  $\mu\text{m}$  film gives an average  $\eta$  value of 2.43% compared with only 2.14% for the SC-3  $\mu\text{m}$  film. The key observation here is that an increase of BOVs in TiO<sub>2</sub> film induces large  $V_{\text{oc}}$  and FF improvements, from 0.797 V and 0.559 for SC-3  $\mu\text{m}$  to 0.826 V and 0.663 for ES-3  $\mu\text{m}$ , respectively. These performance advantages are maintained when a  $\sim 7 \mu\text{m}$  thick TiO<sub>2</sub> film was doctor-bladed on top of the  $\sim 3 \mu\text{m}$  thick electrospun TiO<sub>2</sub> film from a P25 based paste [20], and this film was designated ES-3  $\mu\text{m}$  + P25, as shown in Fig. 6(b). As compared with a 10  $\mu\text{m}$  thick P25 film, the  $V_{\text{oc}}$  and FF values for ES-3  $\mu\text{m}$  + P25 increase by 46 mV and 10.8%, respectively. The cumulative increases of  $V_{\text{oc}}$  and FF give rise to an efficiency of 6.01%, which has increased by 14% as compared with that obtained from a 10  $\mu\text{m}$  thick P25 film.

To investigate the improvement of  $V_{\text{oc}}$  and FF by using ES-8.5 film, charge recombination lifetime ( $\tau_{\text{r}}$ ) [31] was measured through electrochemical impedance spectroscopy (EIS) technique. The impedance at the interface of TiO<sub>2</sub>/dye/electrolyte appeared in the middle frequency range. The bode phase plots (frequency vs. phase angle) of the EIS results for the above two films are shown in Fig. 6(c). The characteristic frequencies at the peaks ( $f_{\text{max}}$ ) is related to charge recombination lifetime ( $\tau_{\text{r}} = 1/2\pi f_{\text{max}}$ ) [32]. As shown in



**Fig. 6.**  $J$ – $V$  curves of DSCs with (a) ES-3  $\mu\text{m}$  and SC-3  $\mu\text{m}$  and (b) ES-3  $\mu\text{m}$  + P25 and P25 film as photoanodes. (c) Electrochemical properties of the bode plot measured by EIS analyzer for DSCs fabricated from ES-3  $\mu\text{m}$ , SC-3  $\mu\text{m}$ , ES-3  $\mu\text{m}$  + P25 and P25 cells. (d) The electron lifetime for ES-3  $\mu\text{m}$  and SC-3  $\mu\text{m}$  cells as a function of  $V_{\text{oc}}$  derived from the open-circuit voltage-decay (OCVD) curve. The corresponding OCVD curves of ES-3  $\mu\text{m}$  and SC-3  $\mu\text{m}$  cells are shown in inset of (d). All the cell areas are 0.45 cm<sup>2</sup>.

Fig. 6(c) and the corresponding parameters in Table 1, the  $\tau_r$  of ES-8.5 (32.2 ms) film is larger than that of SC-8.5 (24.27 ms) film, revealing the increase of the electron lifetime by using TiO<sub>2</sub> film with rich BOVs. This advantage of  $\tau_r$  is maintained when a  $\sim 7 \mu\text{m}$  thick P25 film was doctor-bladed on top of the  $\sim 3 \mu\text{m}$  thick ES-8.5 film. To a certain extent, as electron lifetime increases, recombination current decreases, resulting in an increase of  $V_{oc}$  [8]. FF is usually governed by the series resistance ( $R_s$ ) which can be measured by EIS [33]. However, the  $R_s$  values of ES-8.5 and ES-3  $\mu\text{m}$  + P25 films are very similar to that of SC-8.5 and P25 films, respectively, indicating that the increase of FF for ES-8.5 film is not due to the decrease of  $R_s$ . Further work should be done to explain how BOVs contribute to FF.

The advantages of open-circuit voltage-decay (OCVD) technique over the frequency or steady-state-based methods are that it provides a continuous reading of the lifetime as a function of  $V_{oc}$  at high-voltage resolution. The lifetime can be derived from the open-circuit voltage-decay curve by Equation (1) [34]:

$$\tau_n = -\frac{k_B T}{e} \left( \frac{dV_{oc}}{dt} \right)^{-1} \quad (1)$$

Where,  $k_B T$  is the thermal energy,  $e$  is the elementary charge, and  $dV_{oc}/dt$  is the first-order time derivative of the  $V_{oc}$ . Fig. 6(d) shows lifetime values obtained from the OCVD curves using Equation (1). The shapes of the  $\tau_n$ - $V_{oc}$  curve for ES-8.5 and SC-8.5 shows a trap-assisted conduction mechanism, which are in accordance with that of nanoparticles [6]. The larger  $\tau_n$  for ES-8.5 film compared with that of SC-8.5 may be resulted from the rich BOVs. In addition, the depression seen in the curve at around 0.55 V indicates the effect of SOVs, which could result in recombination of electrons with the electrolyte through tunneling [35]. Therefore, the small difference of the  $\tau_n$  at around 0.55 V for ES-8.5 and SC-8.5 indicates that the two films possess a similar surface texture.

To gain further insight into the effects of the electron traps, we designed a novel cell that has the same structure as the conventional DSCs but without dye adsorption. Without dye molecules, the photovoltaic properties of the cells are just dependent on the inherent electronic structure of TiO<sub>2</sub>. Thus, the effect of BOVs on the photovoltaic properties of DSCs can be enlarged. To find a direct evidence of the effect of BOVs on  $V_{oc}$ , we measured the spectral responses of  $J_{sc}$  and  $V_{oc}$  in the wavelength range from 370 nm to 900 nm for film ES-8.5 and SC-8.5. As shown in Fig. 7, three humps located at 390–450 nm, 450–590 nm and 600–760 nm can be observed for both ES-8.5 and SC-8.5, the locations of which match well with the three Gauss bands fitted from PL spectra. As can be seen from Fig. 7, the  $V_{oc}$  for ES-8.5 film is higher than that of SC-8.5 film across the whole wavelength range from 370 nm to 900 nm; whereas the  $J_{sc}$  is just the reverse. Notably, the peak values of  $V_{oc}$  at 600–760 nm are 0.198 V for SC-8.5 and 0.306 V for ES-8.5, the latter is much higher than the former, which is a direct evidence of the effect of BOVs on the  $V_{oc}$  of DSCs.

Based on the above experiments, we believe that the rich BOVs in ES-8.5 film are the source of the improved  $V_{oc}$  and FF. A schematic diagram was drawn to illustrate the source of the photovoltaic effect related to STE, SOVs and BOVs, as shown in Fig. 8. Generally,

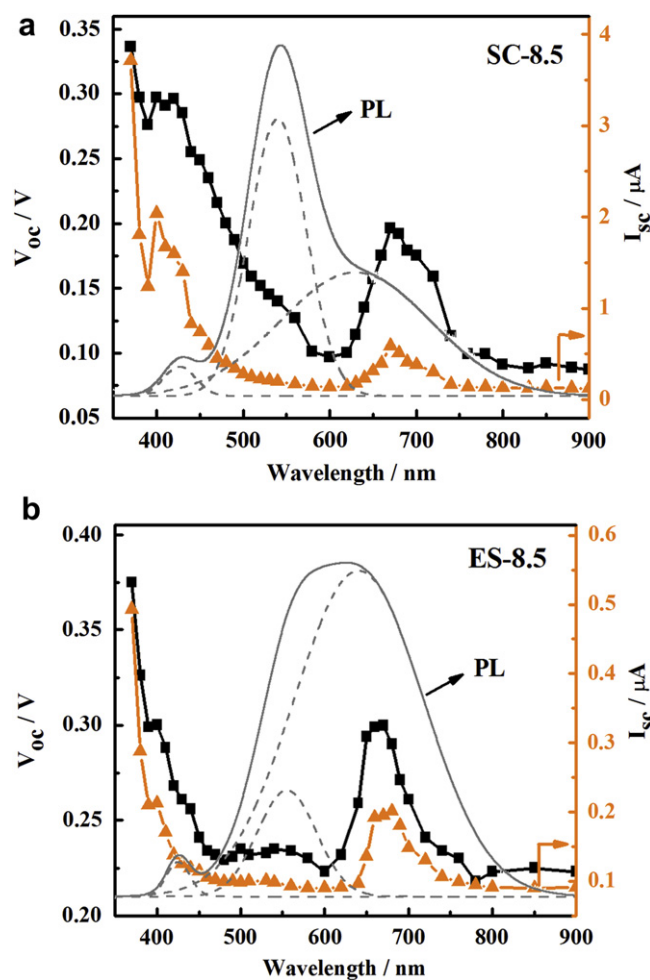


Fig. 7. The spectral responses of the cells with ES-8.5 and SC-8.5 film as photoanodes without dye adsorption. The homogeneous light intensity was fixed at  $200 \text{ mW cm}^{-2}$ .

the steady state DSC performance modeling is solely based on the conduction band electrons and neglects the influence of electron trapping [36]. However, in this work we have demonstrated that the trap-to-trap electron transport and direct recombination from the trap states are indeed happening in DSCs. The electron trapping, transport and recombination processes are shown in Fig. 8. The photons with energy lower than the band gap of TiO<sub>2</sub> promote electrons from the valence band (VB) to trap states in the band gap, followed by trap-to-trap electron transport process and direct recombination from the trap states process. It is necessary to point out that the maximum value of both  $V_{oc}$  and  $J_{sc}$  corresponding to BOVs is at wavelength of 670 nm, which possess a redshift of  $\sim 30 \text{ nm}$  as compared with that of the PL bands from BOVs. This may be because that the traps with energy levels lower than reduction potential of the electrolyte offer an additional pathway to electron transport: the photons promote trapped electrons to trap states in the band gap and/or to CB, followed by trap-to-trap

Table 1

Summary of photovoltaic and electrochemical properties of the cells in Fig. 6. Each value was an average of three measurements.

Samples	$J_{sc}$ [ $\text{mA cm}^{-2}$ ]	$V_{oc}$ [V]	FF	$\eta$ [%]	$R_s$ [ $\Omega$ ]	$\tau_r$ [ms]	Adsorbed dye [ $\times 10^{-8} \text{ mol cm}^{-2}$ ]
ES-8.5	$4.43 \pm 0.10$	$0.826 \pm 0.007$	$0.663 \pm 0.017$	$2.43 \pm 0.06$	$29.83 \pm 1.15$	$32.2 \pm 2.2$	3.79
SC-8.5	$4.81 \pm 0.10$	$0.797 \pm 0.006$	$0.559 \pm 0.007$	$2.14 \pm 0.08$	$29.16 \pm 0.19$	$24.27 \pm 2.67$	3.85
ES-3 $\mu\text{m}$ + P25	$11.96 \pm 0.09$	$0.800 \pm 0.003$	$0.628 \pm 0.003$	$6.01 \pm 0.06$	$26.78 \pm 0.58$	$19.50 \pm 0.33$	12.08
P25	$12.24 \pm 0.11$	$0.750 \pm 0.006$	$0.567 \pm 0.008$	$5.23 \pm 0.04$	$26.57 \pm 0.72$	$17.15 \pm 0.13$	11.85

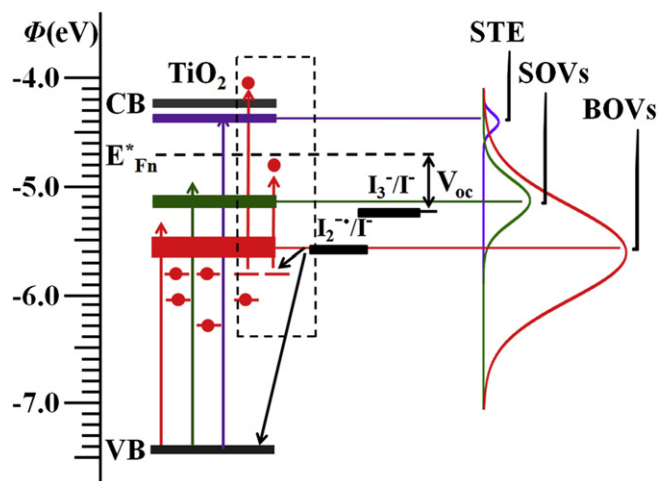


Fig. 8. Schematic illustration showing the source of the photovoltaic effect related to STE, SOVs and BOVs.

electron transport process and/or diffusion process, Fig. 8 (see in rectangle box).

#### 4. Conclusions

In conclusion, with TEOA added to the precursor, transmutation from electrospun nanofibers to transparent nanocrystalline films can be achieved. The proportion of TEOA and the humidity level are two crucial factors in preparation of the transparent nanocrystalline  $\text{TiO}_2$  films. It is found that the emissions from BOVs in electrospun films increase with the proportion of TEOA increasing, and become a dominant peak with 8.5 wt.% TEOA added to the precursor. Contrastively, the dominant peak in PL spectrum of the spin-coated film comes from the emission from SOVs. DSCs fabricated from ES-8.5 film show much higher  $V_{oc}$  and FF than that of SC-8.5. By using a double-layer cell with a doctor-bladed  $\sim 7 \mu\text{m}$  top layer and an electrospun  $\sim 3 \mu\text{m}$  bottom layer, the cumulative increases of  $V_{oc}$  and FF give rise to an efficiency of 6.01%, which has increased by 14% as compared with that obtained from a  $10 \mu\text{m}$  thick P25 film. With the help of EIS measurement and OCVD, we conclude that the BOVs are responsible for the high  $V_{oc}$  and FF.

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#### Appendix A. Supporting information

Supplementary material associated with this article can be found, in the online version, at doi:10.1016/j.jpowsour.2012.04.042

#### References

- [1] B. O'Regan, M. Grätzel, Nature 353 (1991) 737.
- [2] M. Grätzel, Nature 414 (2001) 338.

- [3] M.G. Kang, N.-G. Park, Y.J. Park, K.S. Ryu, S.H. Chang, Sol. Energy Mater. Sol. Cells 75 (2003) 475–479.
- [4] S. Yoon, S. Tak, J. Kim, Y. Jun, K. Kang, J. Park, Build. Environ. 46 (2011) 1899–1904.
- [5] G. Schlichthörl, S.Y. Huang, J. Sprague, A.J. Frank, J. Phys. Chem. B 101 (1997) 8141–8155.
- [6] J. Bisquert, A. Zaban, M. Greenshtein, I. Mora-Seró, J. Am. Chem. Soc. 126 (2004) 13550–13559.
- [7] J. Nelson, S.A. Haque, D.R. Klug, J.R. Durrant, Phys. Rev. B 63 (2001) 205321.
- [8] A. Usami, S. Seki, Y. Mita, H. Kobayashi, H. Miyashiro, N. Terada, Sol. Energy Mater. Sol. Cells 93 (2009) 840–842.
- [9] V. Ganapathy, B. Karunakaran, S.W. Rhee, J. Power Sources 195 (2010) 5138–5143.
- [10] B.A. Gregg, F. Pichot, S. Ferrere, C.L. Fields, J. Phys. Chem. B 105 (2001) 1422–1429.
- [11] K. Park, Q. Zhang, B.B. Garcia, X. Zhou, Y.-H. Jeong, G. Cao, Adv. Mater. 22 (2010) 2329–2332.
- [12] W. Zhang, R. Zhu, X. Liu, B. Liu, S. Ramakrishna, Appl. Phys. Lett. 95 (2009) 043304–043303.
- [13] B. Ding, M. Wang, J. Yu, G. Sun, Sensors 9 (2009) 1609–1624.
- [14] D. Lin, H. Wu, R. Zhang, W. Pan, Chem. Mater. 21 (2009) 3479–3484.
- [15] L. Ji, Y. Yao, O. Toprakci, Z. Lin, Y. Liang, Q. Shi, A.J. Medford, C.R. Millns, X. Zhang, J. Power Sources 195 (2010) 2050–2056.
- [16] H. Wu, L. Hu, M.W. Rowell, D. Kong, J.J. Cha, J.R. McDonough, J. Zhu, Y. Yang, M.D. McGehee, Y. Cui, Nano Lett. 10 (2010) 4242–4248.
- [17] Y.X.D. Li, Adv. Mater. 16 (2004) 1151–1170.
- [18] W. Wang, J. Zhou, S. Zhang, J. Song, H. Duan, M. Zhou, C. Gong, Z. Bao, B. Lu, X. Li, W. Lan, E. Xie, J. Mater. Chem. 20 (2010) 9068–9072.
- [19] Y. Su, B. Lu, Y. Xie, Z. Ma, L. Liu, H. Zhao, J. Zhang, H. Duan, H. Zhang, J. Li, Y. Xiong, E. Xie, Nanotechnology 22 (2011) 285609.
- [20] X. Li, Y. Zhang, Z. Zhang, J. Zhou, J. Song, B. Lu, E. Xie, W. Lan, J. Power Sources 196 (2011) 1639.
- [21] R. Zhu, C.-Y. Jiang, X.-Z. Liu, B. Liu, A. Kumar, S. Ramakrishna, Appl. Phys. Lett. 93 (2008) 013102–013103.
- [22] T. Sugimoto, T. Kojima, J. Phys. Chem. C 112 (2008) 18760–18771.
- [23] H. Tang, H. Berger, P.E. Schmid, F. Lévy, G. Burri, Solid State Commun. 87 (1993) 847–850.
- [24] Y. Lei, L.D. Zhang, G.W. Meng, G.H. Li, X.Y. Zhang, C.H. Liang, W. Chen, S.X. Wang, Appl. Phys. Lett. 78 (2001) 1125–1127.
- [25] N. Serpone, D. Lawless, R. Khairutdinov, J. Phys. Chem. 99 (1995) 16646–16654.
- [26] L.V. Saraf, S.I. Patil, S.B. Ogale, S.R. Sainkar, S.T. Kshirsager, Int. J. Mod. Phys. B 12 (1998) 2635–2647.
- [27] L. Forss, M. Schubnell, Appl. Phys. B 56 (1993) 363–366.
- [28] D.C. Cronemeyer, Phys. Rev. 113 (1959) 1222.
- [29] F. Montoncello, M.C. Carotta, B. Cavicchi, M. Ferroni, A. Giberti, V. Guidi, C. Malagu, G. Martinelli, F. Meinardi, J. Appl. Phys. 94 (2003) 1501–1505.
- [30] A. Amout, R. Leonelli, Phys. Rev. B 51 (1995) 6842.
- [31] Q. Wang, J.-E. Moser, M. Grätzel, J. Phys. Chem. B 109 (2005) 14945–14953.
- [32] J. Bisquert, J. Phys. Chem. B 106 (2001) 325–333.
- [33] G. Kron, U. Rau, J.H. Werner, J. Phys. Chem. B 107 (2003) 13258–13261.
- [34] A. Zaban, M. Greenshtein, J. Bisquert, ChemPhysChem 4 (2003) 859.
- [35] P.S. Archana, R. Jose, C. Vijila, S. Ramakrishna, J. Phys. Chem. C 113 (2009) 21538–21542.
- [36] J. Halme, P. Vahermaa, K. Miettunen, P. Lund, Adv. Mater. 22 (2010) E210–E234.

#### A list of abbreviations

- BOVs: bulk oxygen vacancies  
 SOVs: surface oxygen vacancies  
 STE: self-trapped exciton  
 OV: oxygen vacancies  
 PL: photoluminescence  
 DSCs: dye-sensitized solar cells  
 EIS: electrochemical impedance spectroscopy  
 OCVD: open-circuit voltage-decay  
 TEOA: tri-ethanolamine  
 ES: electrospinning  
 SC: spin-coating  
 ES-8.5: film prepared by electrospinning with 8.5 wt.% TEOA  
 SC-8.5: film prepared by spin-coating with 8.5 wt.% TEOA  
 ES-3  $\mu\text{m}$ : 3  $\mu\text{m}$  thick ES-8.5 film  
 SC-3  $\mu\text{m}$ : 3  $\mu\text{m}$  thick SC-8.5 film  
 ES-3  $\mu\text{m}$  + P25: a double-layer film of  $\sim 7 \mu\text{m}$  thick P25 film on top of the  $\sim 3 \mu\text{m}$  ES-8.5 film  
 P25: 10  $\mu\text{m}$  thick doctor-bladed P25 film